

# Spontaneous symmetry breaking and finite-size spectra of quantum frustrated antiferromagnets

Proceedings: [cond-mat/0607764](#)



Grégoire Misguich

Service de Physique Théorique (SPhT)

CEA Saclay, France

Philippe Sindzingre

LPTMC, Paris-6 University, France

+ many thanks to:

B. Bernu, S. Furukawa, A. Läuchli, C. Lhuillier,  
M. Oshikawa, V. Pasquier, K. Penc, L. Pierre & N. Shannon

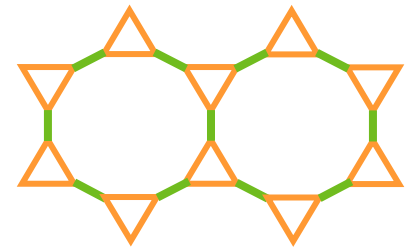
# Outline

---

- Exact diagonalizations and quantum numbers
- Spontaneous symmetry breaking
- Discrete SSB: some examples of valence-bond crystals

Kagome lattice, expanded kagome lattice, others...

- Continuous SSB: Néel and nematic phases



# Exact diagonalizations and quantum numbers

- Quantum lattice model  $\Rightarrow$  symmetry group  $G$
- $G \Rightarrow$  Irreducible representations (irreps)

Possible “automatic” implementation. Using GAP/GRAPe for example ([www.gap-system.org](http://www.gap-system.org))

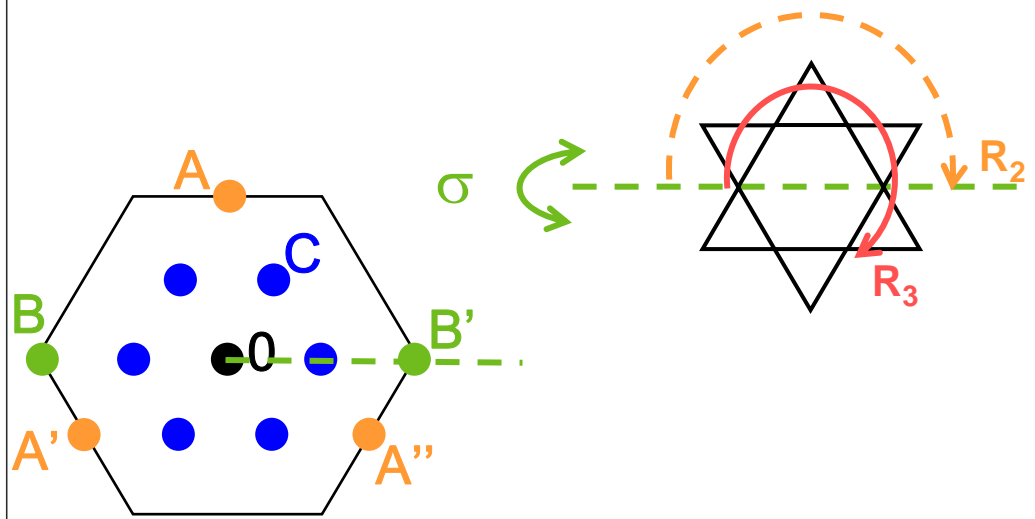
- Irrep.  $\Rightarrow$  block diagonalization

In most practical cases, irreps are induced by *1d representations* of some subgroup.

First fix momentum  $\mathbf{k}$ , then look for irrep. of the “little” (point) group of  $\mathbf{k}$

- Example. Irreps of the first levels of a 36-site kagome Heisenberg model

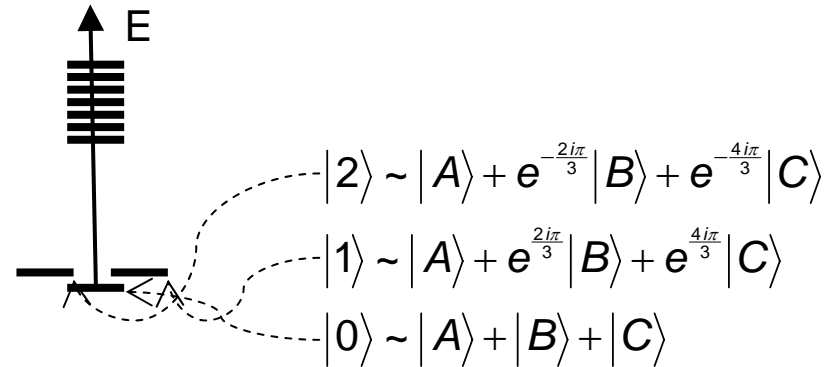
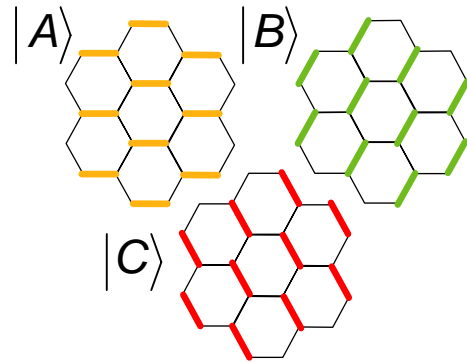
Number	$2\langle \vec{S}_i \cdot \vec{S}_j \rangle$	$k$	$R_3$	$R_2$	$\sigma$	Deg.
1	-.43837653	0	1	1	1	1
2	-.43809562	$B$	$e^{\pm 2i\pi/3}$			4
3	-.43807091	0	$e^{\pm 2i\pi/3}$	1		2
4	-.43799346	0	1	1	1	1
5	-.43785105	$C$			1	6
6	-.43758510	0	1	-1	1	1
7	-.43758455	$A$		1	1	3
8	-.43751941	$C$			-1	6
9	-.43721566	0	1	1	-1	1
10	-.43718796	0	$e^{\pm 2i\pi/3}$	1		2
11	-.43714765	$A$		-1	-1	3
12	-.43705108	0	$e^{\pm 2i\pi/3}$	-1		2
13	-.43703981	$B$	1		1	2
14	-.43703469	$A$		-1	1	3
15	-.43685867	0	1	-1	-1	1
16	-.43685319	$B$	1		-1	2
17	-.43683757	$A$		1	-1	3



[Waldtmann, EPJB 1998 & Sindzingre, unpublished]

# Spontaneous symmetry breaking and quantum numbers

- Broken symmetry states  $\neq$  eigenstates



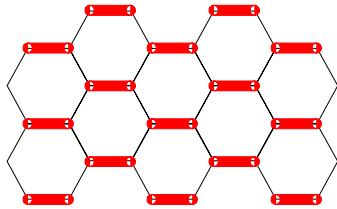
- SSB  $\Rightarrow$  quasi-degenerate ground-state in finite-size system, *several* irreps.
- Group theory  $\Rightarrow$  Predict the irreps associated to a given SSB phase

$$n_\gamma = \frac{1}{|G|} \sum_{g \in G} \underbrace{\chi_\gamma(g^{-1})}_{\text{character}} \sum_{i=1}^d \langle i | \hat{g} | i \rangle$$

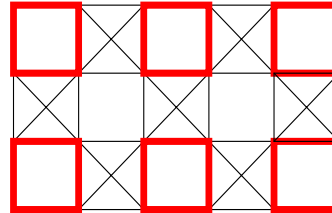
|A>, ... : Basis of the ground-state subspace

- A simple (no fluctuations) trial ground-state can be used to compute  $n_\gamma$
- Simple for *discrete* SSB, slightly more involved for *continuous* ones

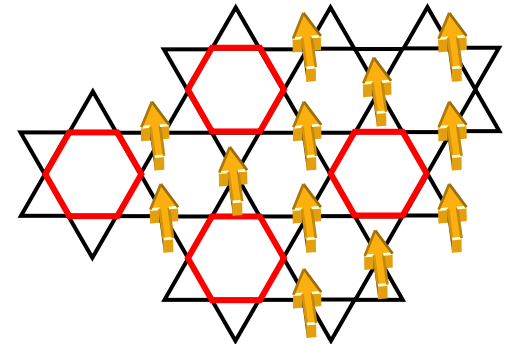
# Examples of valence-bond crystals (from ED studies)



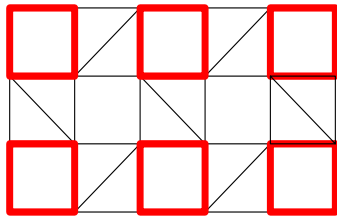
$J_1$ - $J_2$ - $J_3$  model  
Fouet *et al.* EPJB 2001



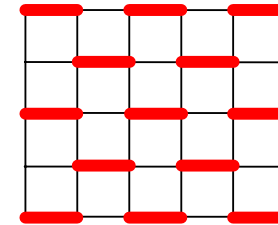
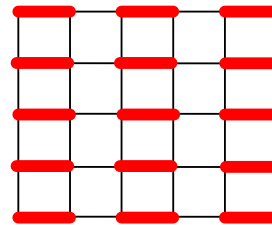
Fouet *et al.* PRB 2003



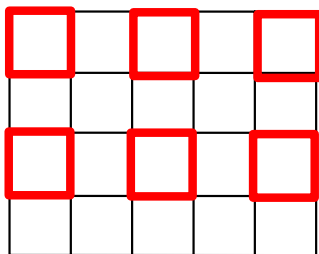
Kagome, 1/3 magnetization plateau  
Cabra *et al.*, PRB 2005



Shastry-Sutherland lattice  
Läuchli, Wessel & Sigrist PRB 2002



Heisenberg model & 4-spin "ring" exchange  
Läuchli *et al.* PRL 2005

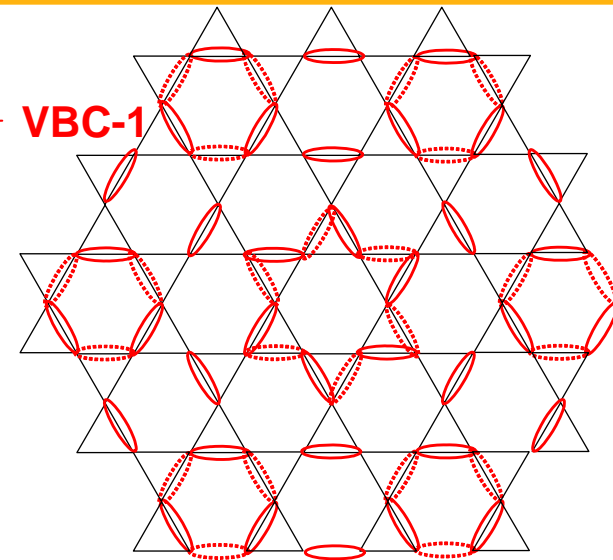


$J_1$ - $J_2$ - $J_3$  model  
Mambrini *et al.*, cond-mat/0606776

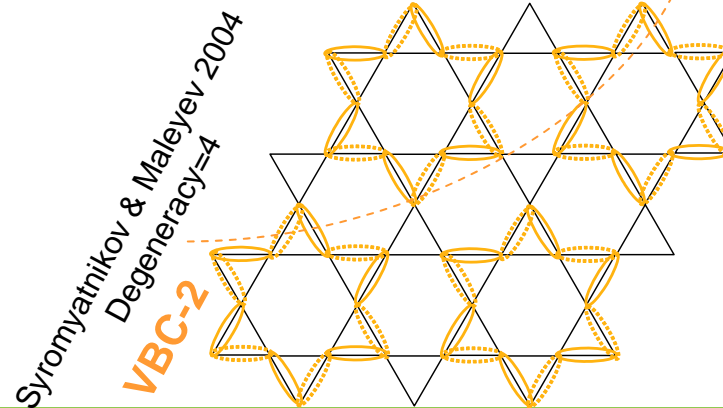
+ others...

# Valence bond crystals on kagome ?

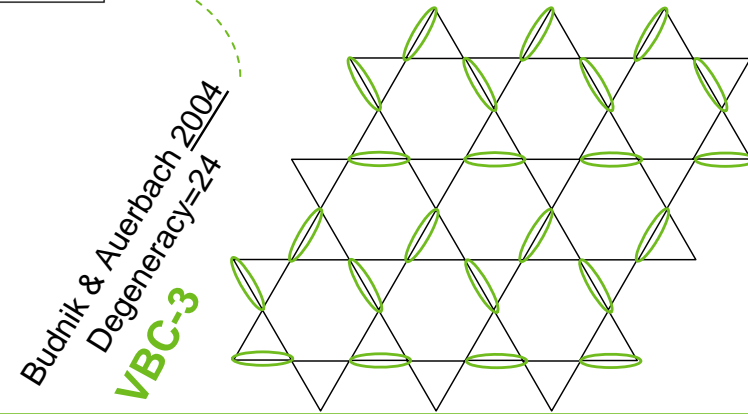
Number	$2\langle \vec{S}_i \cdot \vec{S}_j \rangle$	$k$	$R_3$	$R_2$	$\sigma$	Deg.	VBC-1	VBC-2	VBC-3
1	-0.43837653	0	1	1	1	1		*	*
2	-0.43809562	$B$	$e^{\pm 2i\pi/3}$			4			*
3	-0.43807091	0	$e^{\pm 2i\pi/3}$	1		2			*
4	-0.43799346	0	1	1	1	1			*
5	-0.43785105	$C$			1	6			*
6	-0.43758510	0	1	-1	1	1	*		*
7	-0.43758455	$A$		1	1	3	*		*
8	-0.43751941	$C$			-1	6	*		*
9	-0.43721566	0	1	1	-1	1	*		*
10	-0.43718796	0	$e^{\pm 2i\pi/3}$	1		2			*
11	-0.43714765	$A$		-1	-1	3			*
12	-0.43705108	0	$e^{\pm 2i\pi/3}$	-1		2			*
13	-0.43703981	$B$	1		1	2			*
14	-0.43703469	$A$		-1	1	3			*
15	-0.43685867	0	1	-1	-1	1			*
16	-0.43685319	$B$	1		-1	2			*
17	-0.43683757	$A$		1	-1	3		*	*
...	...	...	...	...	...	...			*
44	-0.43474519	0	$e^{\pm 2i\pi/3}$	-1		2			*



Marston & Zeng 1991; Nikolic & Senthil 2003  
Degeneracy=12



Syromyatnikov & Maleyev 2004  
Degeneracy=4  
**VBC-2**

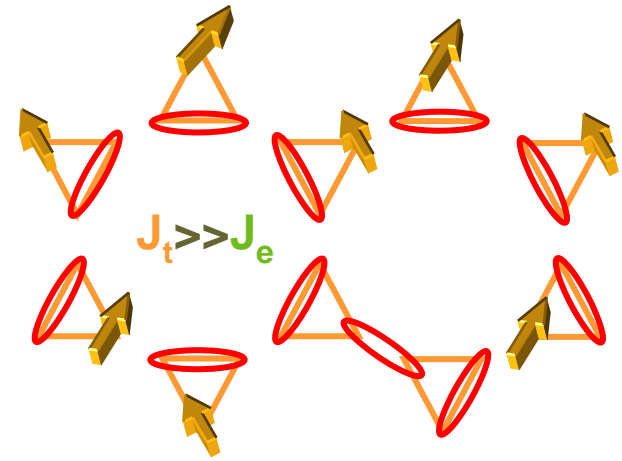
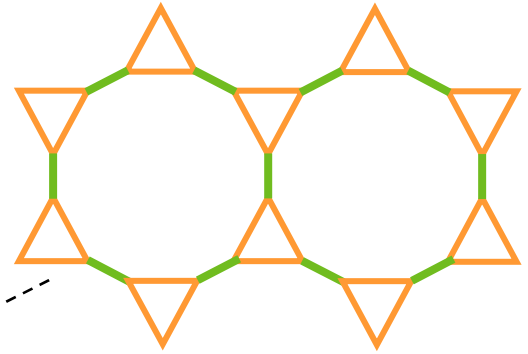
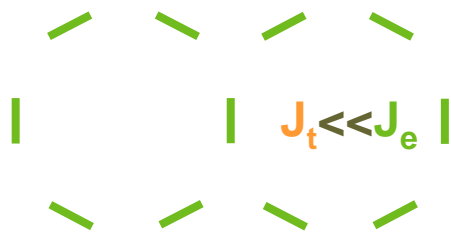


Budnik & Auerbach 2004  
Degeneracy=24  
**VBC-3**

# Expanded kagome (star) lattice $J_e$ - $J_t$ model

Similarities with kagome:

- same classical AF degeneracy
- same # of dimer coverings

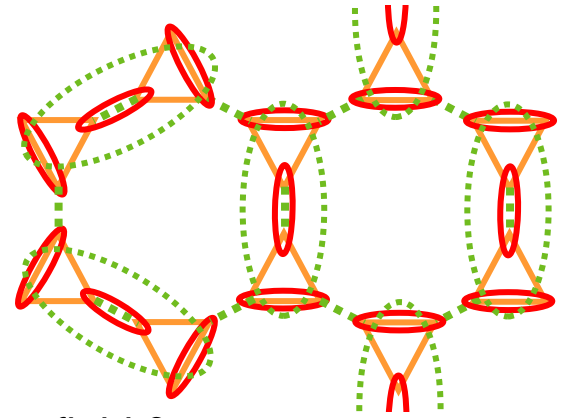


Isolated dimers - trivial limit  
 Single singlet ground-state  
 Large gap to all excitations  
 “quantum paramagnet”

At  $J_t=J_e$ , the model is in this phase  
 [Richter et al. [PRB 2004](#)]

Extensive degeneracy  
 Degenerate perturbation theory  
 Effective spin-chirality Hamiltonian [Subrahmanyam [PRB 1995](#)]

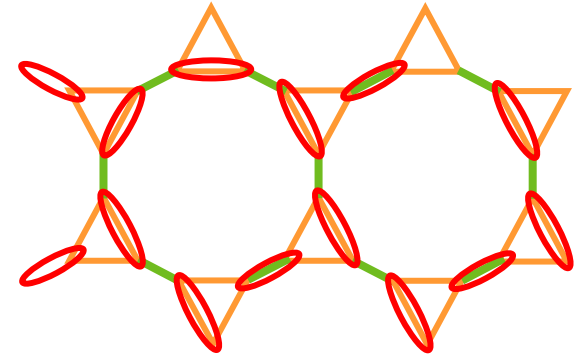
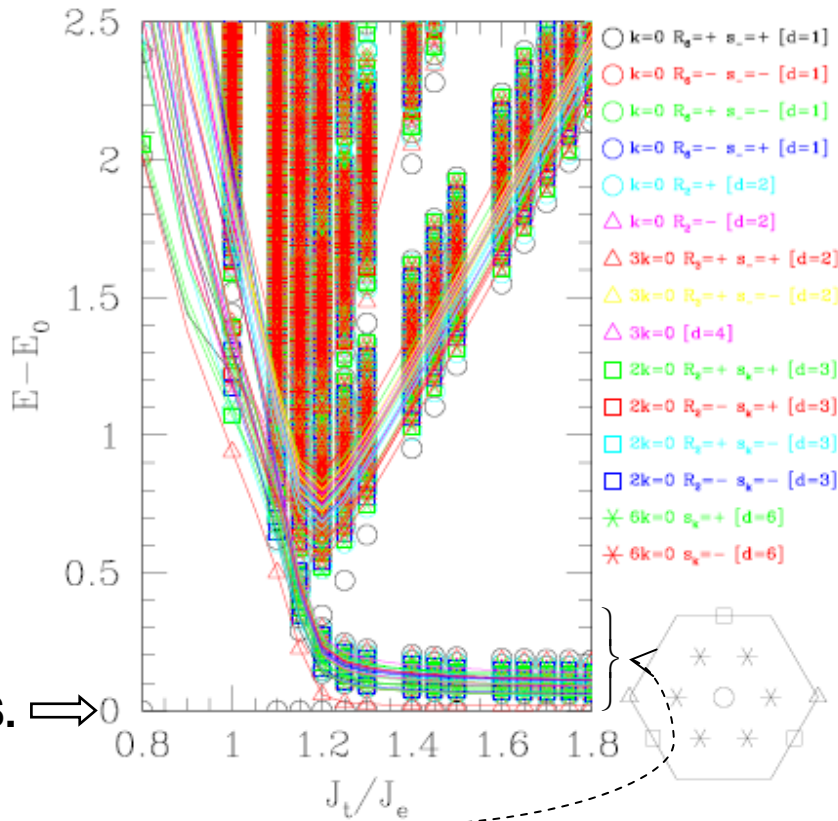
Mean-field approximation  
 Degenerate solutions = super coverings  
 [same as in Mila [PRL 1998](#) !]



How is this degeneracy lifted beyond mean-field ?  
 ⇒ We tried ED in the 1<sup>st</sup> neighbor valence-bond subspace

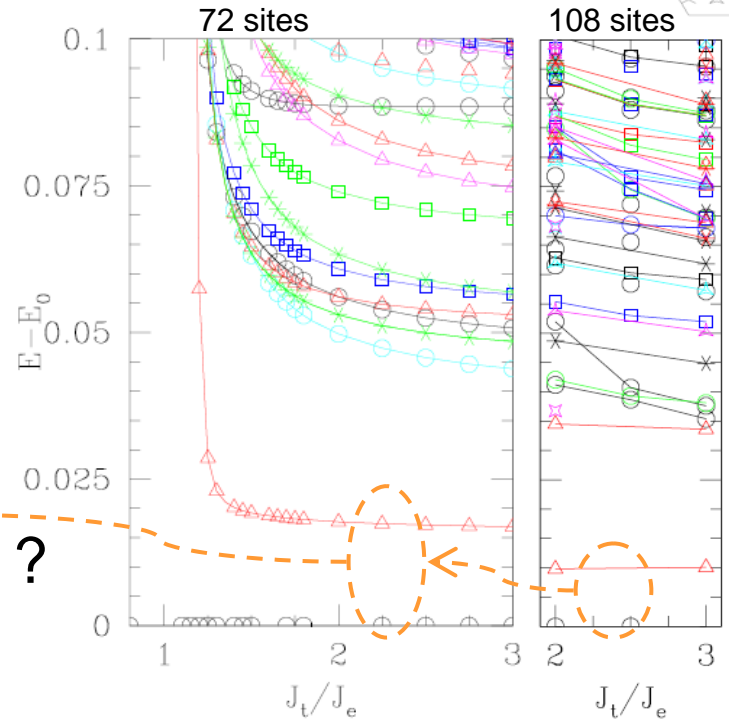
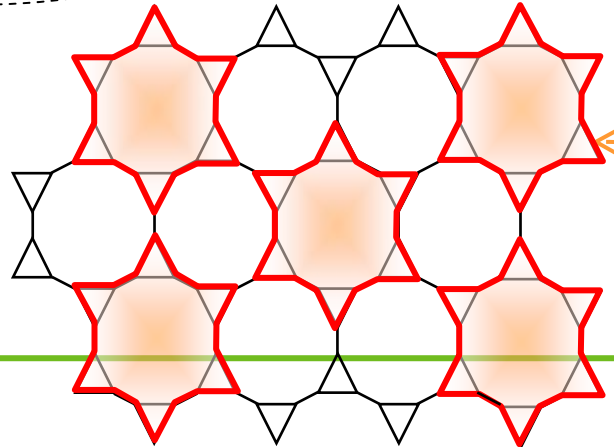
# Expanded kagome lattice $J_e$ - $J_t$ model: a VBC when $J_t \gg J_e$ ?

ED in the 1<sup>st</sup> neighbor valence-bond subspace  
 [Zeng & Elser *PRB* 1995; Marnbrini & Mila *EPJB* 2000]



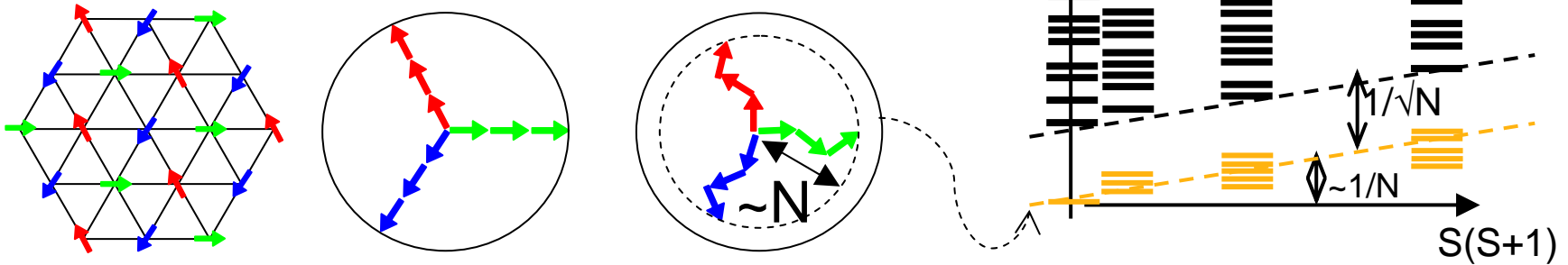
G.S.  $\Rightarrow$

# of super coverings  
(mean-field) states





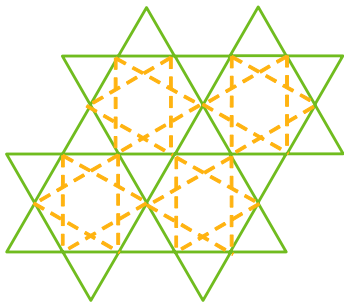
# Néel phases – Anderson tower



Anderson, PR 1953

Bernu *et al.*, PRL 1992, PRB 1994

Lhuillier, [cond-mat/0502464](https://arxiv.org/abs/cond-mat/0502464)

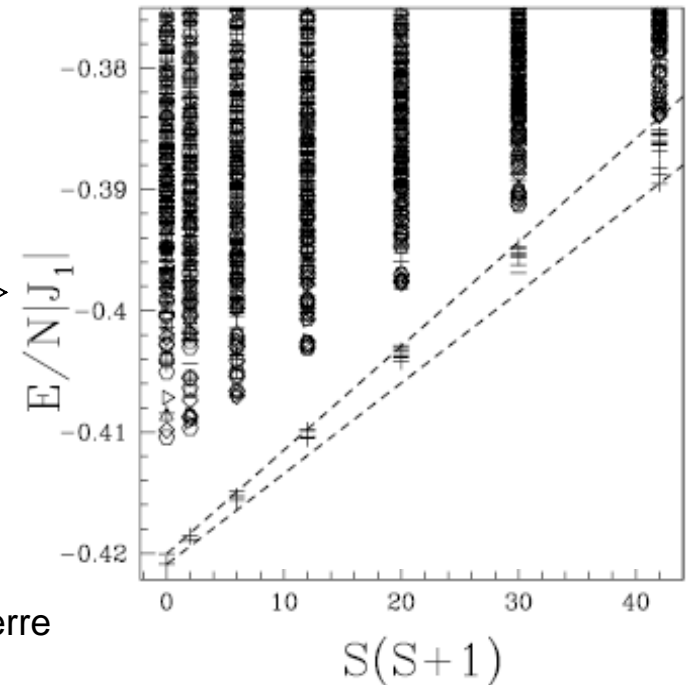
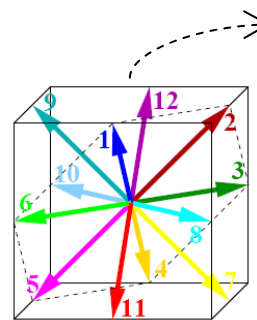


Example with 12 sublattices !

$J_1 - J_2$  model on the kagome lattice

J.-C. Domenge, P. Sindzingre, C. Lhuillier and L. Pierre

PRB 2005



# Quantum spin nematics

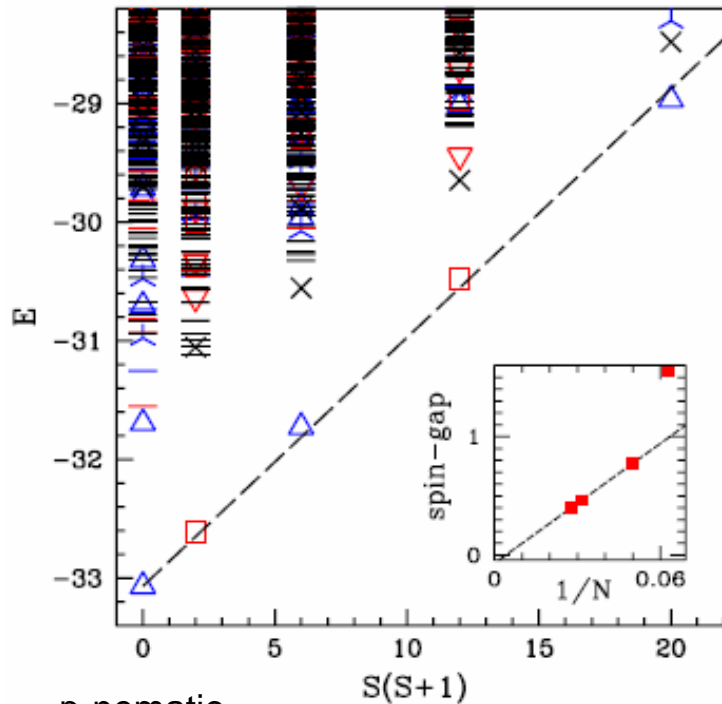
[Andreev & Grishchuk JETP 1984; Chandra & Coleman PRL 1991]

- Spontaneous breakdown of SU(2) symmetry but short-ranged spin-spin correlations. The order parameter involves 2 spins operators

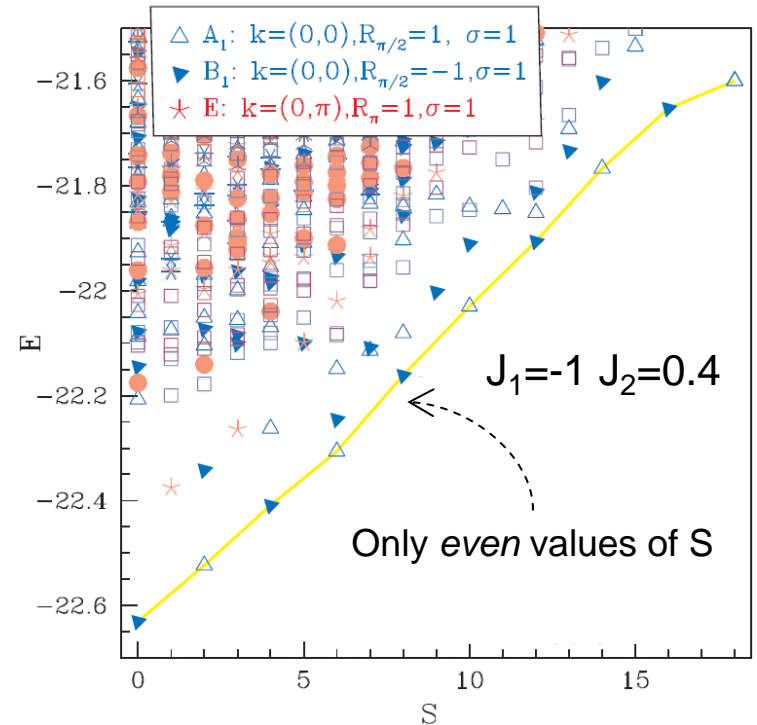
$$\langle \vec{S}_r \rangle = 0$$

Spontaneous selection of a (oriented or non oriented) *plane*

$N=32$   $\theta/\pi \approx 0.31$  ( $K/J=1.5$ )



p-nematic  
Spin-  $\frac{1}{2}$  model with 4-spin ring exchange (square lattice)  
Läuchli *et al.*, [PRL 2005](#)



n-nematic  
Shannon, Momoi & Sindzingre, [PRL 2006](#)

# Conclusions

---

- Exact diagonalizations, when coupled to a full symmetry analysis is a powerful tool specially for  $D > 1$  quantum frustrated magnets
- The spectrum itself is a rich source of information concerning broken symmetries
- Very useful if only very small systems are numerically available  
(spin  $> \frac{1}{2}$ , spin-orbital models, doped magnets, ...)
- Complementary approach to obtain the order parameter from the quasi-degenerate wave-functions (using reduced density matrices): Furukawa, GM & Oshikawa [PRL 2006](#)