



School and Workshop on
Highly Frustrated Magnets and
Strongly Correlated Systems:
From Non-Perturbative Approaches to
Experiments

30 July - 17 August 2007

Miramare-Trieste, Italy



2D quantum magnetism and spin liquids



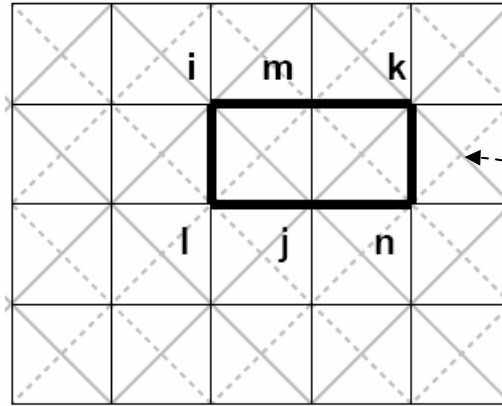
Grégoire Misguich

Service de Physique Théorique
Commissariat à l'Énergie Atomique (CEA)
Centre de Saclay, France

www-spht.cea.fr/pisp/misguich

Tutorial: Exact VBC ground-states in a toy model

Gellé, Lauchli, Kumar & Mila, [arxiv:0704.2352v1](https://arxiv.org/abs/0704.2352v1)



Sum over all 6-site rectangles

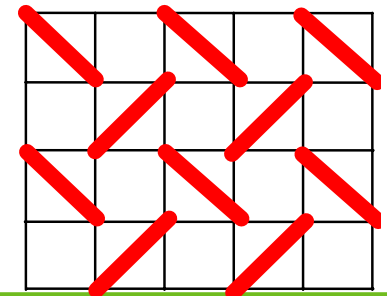
$$P_{i,j,k}^A = |\vec{S}_i + \vec{S}_j + \vec{S}_k|^2 - \frac{3}{4}$$

$$P_{l,m,n}^B = |\vec{S}_l + \vec{S}_m + \vec{S}_n|^2 - \frac{3}{4}$$

Projectors onto the states with total spin $S=3/2$

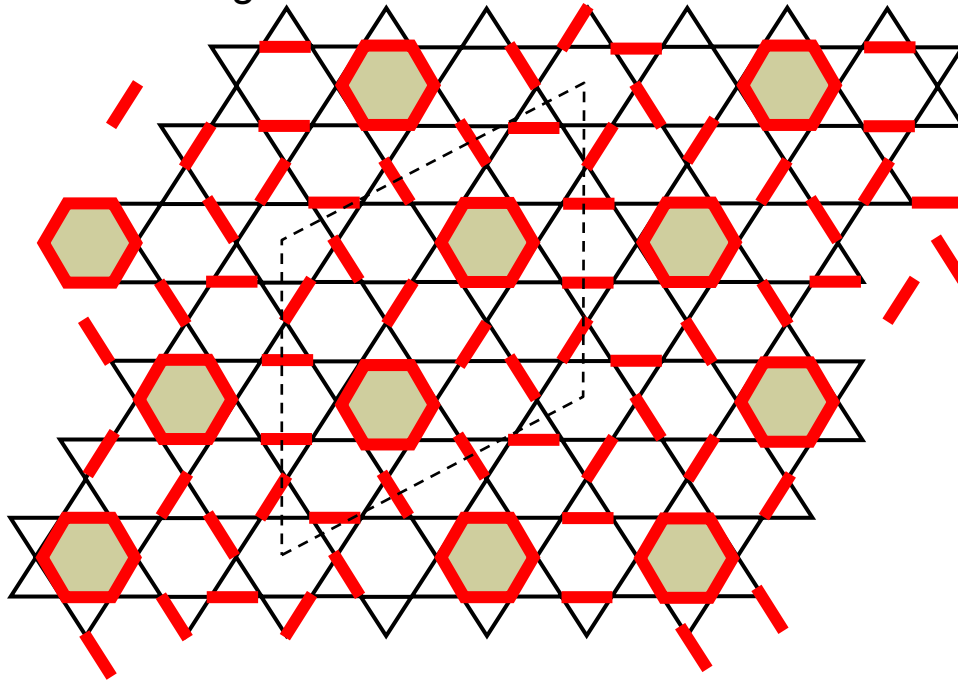
$$H_0 = \sum_{[i,j,k,l,m,n]} \frac{1}{4} P_{i,j,k}^A P_{l,m,n}^B$$

□ Show that this singlet-product state :
is an exact ground-state of H_0 :



Tutorial – VBC stabilization by local resonances

- ❑ What is the ground-state degeneracy of the VBC below ? How many sites and bonds per unit cell ?
- ❑ Compare the variational energy of some arbitrary first-neighbor valence bond covering on the kagome lattice with that of the VBC below (with resonance around each shaded hexagon)
Hint : the ground-state energy of a periodic Heisenberg chain with 6 sites is $E=-2.803$

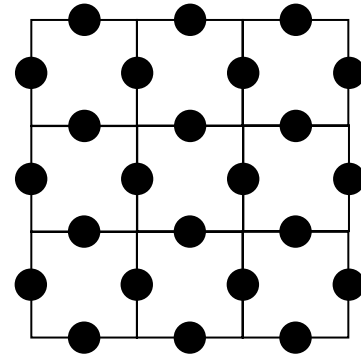


Such VBC phase has been proposed for the spin- $\frac{1}{2}$ Heisenberg model on the kagome lattice: see Marston & Zeng [1991](#); Nikolic & Senthil [2003](#); Singh & Huse [2007](#)

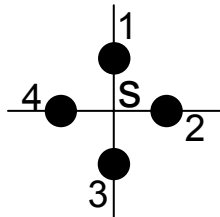
Tutorial – Kitaev’s model

Simplest spin model with (Z_2) fractionalization.

Elementary excitations are similar to that of a short-range RVB liquid
 Kitaev arXiv 1997 (Annals of Phys. [2003](#)).



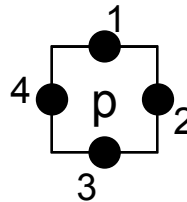
- Spin-1/2 model where the spins leaves on the bonds



$$H = -\sum_s A_s - \sum_p B_p$$

$$A_s = \sigma_1^x \sigma_2^x \sigma_3^x \sigma_4^x$$

$$B_p = \sigma_1^z \sigma_2^z \sigma_3^z \sigma_4^z$$



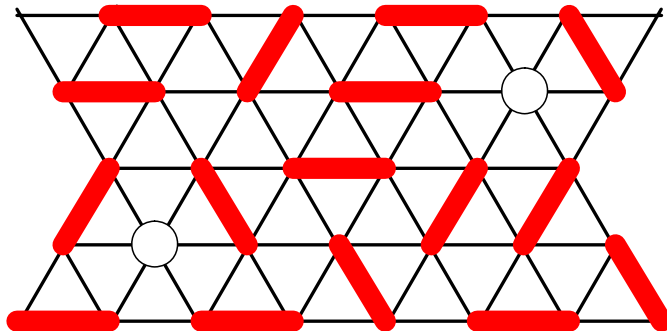
- Show that $[A_s, B_p]=0$
- Consider the equal amplitude superposition $|0\rangle$ of all the spin configurations (in the σ^z basis) with an *even* number of up down spins on each plaquette p .
 Show that $|0\rangle$ is a ground state of H
- We call “vison” a plaquette p with $B_p=-1$. Show that the number of “vison” is always even.
- We call “spinon” a vertex s with $A_s=-1$. Show that the number of “spinon” is always even.
- Construct some eigenstate of H with two remote “visons”. Do the same with two “spinons”.
- Show that the ground-state is two-fold degenerate on a cylinder. Hint: consider a product of σ^z operators on a loop winding around the cylinder.

Tutorial: non-magnetic impurities as probes for spinon (de)confinement

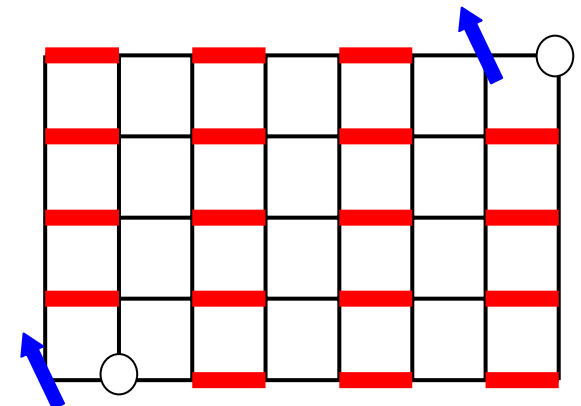
Sachdev & Vojta, [cond-mat/0009202](https://arxiv.org/abs/cond-mat/0009202)

- Consider two spin systems: one in a columnar VBC phase, and the other in a short-ranged RVB phase.
- We remove two spins (non-magnetic “impurities”) at some fixed large distance from each other.
- In both cases, what are the lowest energy states (ground-state degeneracy ? Total spin ?)

□ Deconfined short-range RVB phase



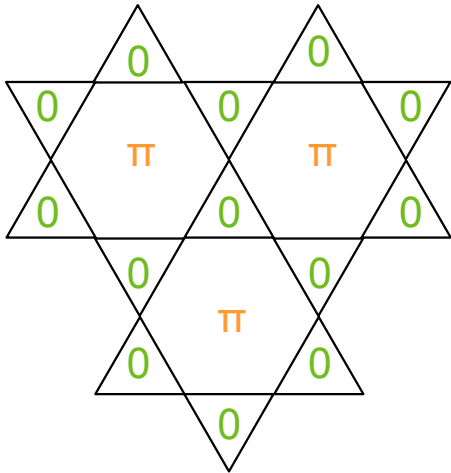
□ Valence-bond crystal



Tutorial: a mean-field spin-liquid state on the kagome lattice

M. B. Hastings, Phys. Rev. B 63, 014413 (2000)

Y. Ran, M. Hermele, P. Lee & X.-G. Wen, Phys. Rev. Lett. 98, 117205 (2007)



$$\frac{J_{ij}}{8} \begin{bmatrix} \langle c_{i\uparrow} c_{j\uparrow} + c_{i\downarrow} c_{j\downarrow} \rangle & \langle c_{i\uparrow} c_{j\downarrow} - c_{i\downarrow} c_{j\uparrow} \rangle \\ \langle c_{i\downarrow} c_{j\uparrow} - c_{i\uparrow} c_{j\downarrow} \rangle & \langle c_{i\downarrow} c_{j\downarrow} + c_{i\uparrow} c_{j\uparrow} \rangle \end{bmatrix} = \begin{bmatrix} -\chi_{ij}^+ & \eta_{ij} = 0 \\ \eta_{ij}^+ = 0 & \chi_{ij} \end{bmatrix}$$

$$H_{MF} = \sum_{\langle ij \rangle} \chi_{ij} (c_{i\downarrow}^+ c_{j\downarrow} + c_{i\uparrow}^+ c_{j\uparrow}) + H.c$$

$$\chi_{ij} = \pm 0.221 \quad \prod_{\text{triangle}} \chi_{ij} = 1 \quad \prod_{\text{hexagon}} \chi_{ij} = -1$$

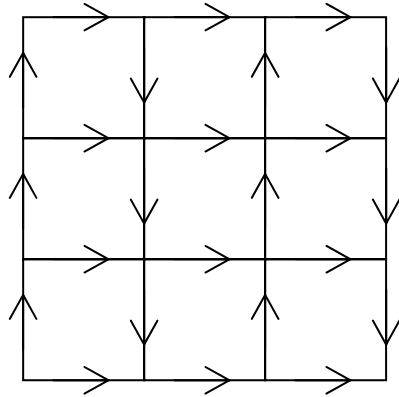
- Find some χ_{ij} such that the flux is π per hexagon, and 0 per triangle.
- Show that the invariant gauge group is U(1).

Tutorial: π -flux state

π -flux state

Affleck & Marston, Phys. Rev. B 37, 3774 (1988)

(equivalent to the “mixed” $s + id$ RVB state of Kotliar & Liu, Phys. Rev. B 38, 5142 (1988))



$$U_{i \rightarrow j}^0 = i \begin{bmatrix} u & 0 \\ 0 & u \end{bmatrix} \quad u \in \mathfrak{R}$$

$$\vec{A}_i^0 = 0$$

$$H_{MF} = 2u \sum_{\langle i \rightarrow j \rangle} (i c_{\uparrow i}^+ c_{\uparrow j} + \text{H.c.})$$

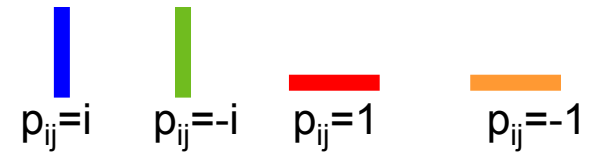
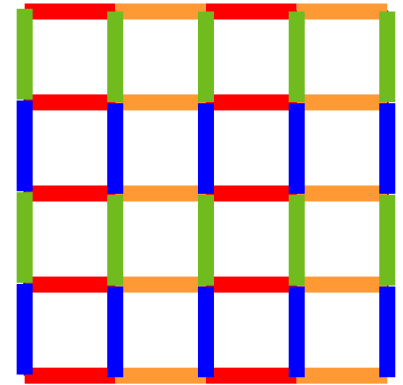
- Compute the SU(2) “flux” $U_{12} U_{23} U_{34} U_{41}$ piercing each square plaquette.
- Show that this state is translation invariant
- Show that invariant gauge group of this state is SU(2)
- Show that the spinon dispersion relation is:

$$E_k = \pm 4u \sqrt{\sin^2(k_x) + \sin^2(k_y)}$$

Tutorial: valence-bond crystal order parameters

- Show that P is an order parameter for a columnar VBC

$$\begin{aligned}
 P &= \frac{1}{N} \langle 0 | \left| \sum_{\langle ij \rangle} p_{ij} \vec{S}_i \cdot \vec{S}_j \right|^2 | 0 \rangle \\
 &= \frac{1}{N} \sum_{\langle ij \rangle} \sum_{\langle kl \rangle} p_{ij} \bar{p}_{kl} \underbrace{\langle 0 | (\vec{S}_i \cdot \vec{S}_j) (\vec{S}_k \cdot \vec{S}_l) | 0 \rangle}_{\text{4-spin correlation}}
 \end{aligned}$$



- Find some order parameter(s) which can distinguish a plaquette and a columnar phases:

